

# Potentially Relevant Techniques from QCD/Lattice

Lance Dixon

Physics Opportunities at a Lepton Collider  
in the Fully Nonperturbative QED Regime

SLAC

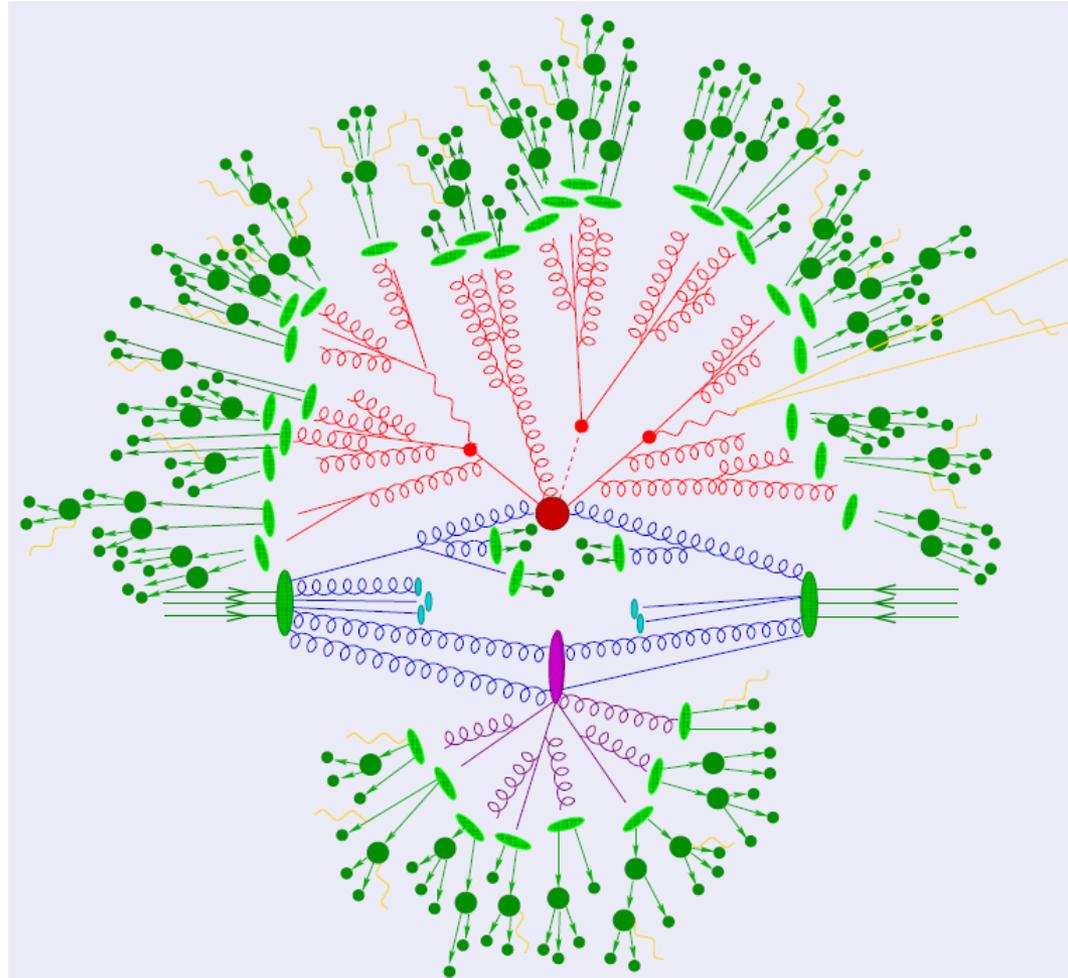
August 8, 2019

# Outline

- Lessons from perturbative QED?
- Comments on Ritus-Narozhny conjecture
- Lessons from lattice QCD/QED?

# Lessons from perturbative QCD, e.g. for LHC?

- Seems like a very different setup:  
Premium on many external particles,  
no external fields.  
Quarks and gluons massless.

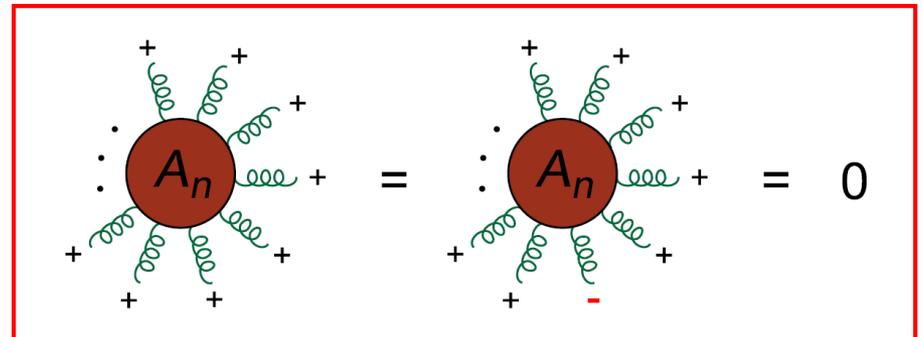
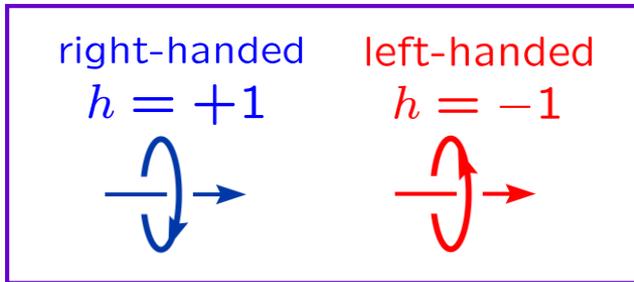


# Interesting all- $n$ helicity properties

Grisaru, Pendleton, van Nieuwenhuizen (1977)

Many **tree-level** multi-gluon **helicity** amplitudes either vanish or are very short

All outgoing helicity convention



Identities derived using supersymmetry (superpartners don't contribute!).

Can also **replace 2 gluons by massless quarks**

$$A_n = \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}$$

Parke-Taylor formula (1986)

# QED application

Similarly,

$$= 0 \quad \text{for } m_e \rightarrow 0$$

- Holds for any photon and electron momenta.
- Now let all but one of the photon momenta be the **same**, i.e. come from **same** circularly polarized laser beam. Implies:

$$= 0 = \quad \text{for } m_e \rightarrow 0$$

# Some questions

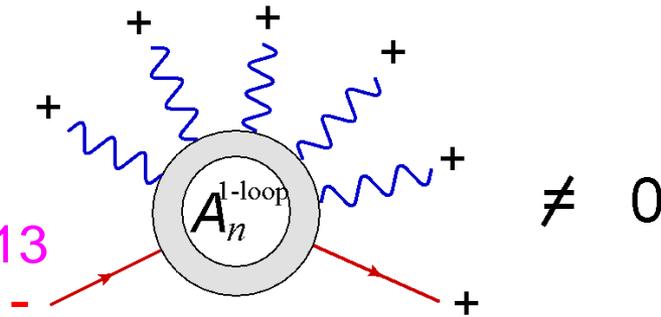
- Is this (semi)-strong field relation known already?
- Can it be tested experimentally?  
 (“Test of supersymmetry”?)
- With optical lasers (1eV, 1000nm) need  $E_e > 500$  GeV to get  $s_{e\gamma} \gg m_e^2$
- With 10nm bunch lepton collider, should be in this regime, but circular polarization??
- Focused LCLS beam (1keV, 1nm) [talk by Claudio Pellegrini] only needs  $E_e > 1$  GeV. Circular polarization? How to measure outgoing  $\gamma$  polarization?

# Quantum corrections?

- Amplitudes that vanish at tree level due to supersymmetry **generically** become nonvanishing at one-loop (although they have vanishing branch cuts).

- Indeed,

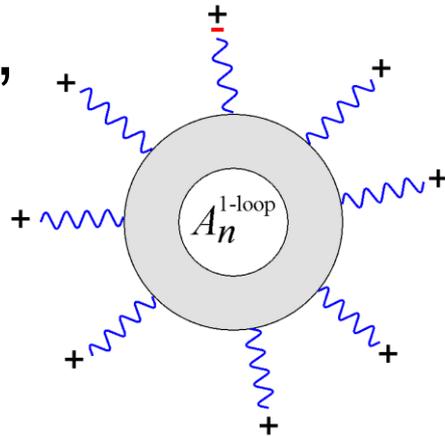
G. Mahlon, 9311213



- However,

also

Bern, Chalmers,  
LD, Kosower,  
hep-ph/9312333  
and BDK,  
hep-ph/0505055



both for  $m_e \rightarrow 0$

= 0

except for  $n = 4$

cf. Dinu, Heinzl, Ilderton,  
Marklund, Torgrimsson,  
1312.6419 for  $b_0 \gg 1$ ?

# Lessons from strong-field QED for QCD?

Gluon helicity flip in a plane wave background

1903.01491

---

Tim Adamo<sup>1</sup> & Anton Ilderton<sup>2</sup>

<sup>1</sup>*Theoretical Physics Group, Blackett Laboratory  
Imperial College London, SW7 2AZ, United Kingdom*

<sup>2</sup>*Centre for Mathematical Sciences  
University of Plymouth, PL4 8AA, United Kingdom*

*E-mail: t.adamo@imperial.ac.uk, anton.ilderton@plymouth.ac.uk*

ABSTRACT: We compute the leading probability for a gluon to flip helicity state upon traversing a background plane wave gauge field in pure Yang-Mills theory and QCD, with an arbitrary number of colours and flavours. This is a one-loop calculation in perturbative gauge theory around the gluonic plane wave background, which is treated without approximation (i.e., to all orders in the coupling). We introduce a background-dressed version of the spinor helicity formalism and use it to obtain simple formulae for the flip amplitude with pure external gluon polarizations. We also give in-depth examples for gauge group  $SU(2)$ , and evaluate both the high- and low-energy limits. Throughout, we compare and contrast with the calculation of photon helicity flip in strong-field QED.

# Comment on Ritus-Narozhny conjecture

Fedotov, 1608.02261, and talk yesterday

Let  $\chi = \alpha^{-3/2}$

$$\begin{aligned}
 \frac{M}{m} = & \underbrace{\text{[Diagram 1]}}_{\simeq \alpha \chi^{2/3} \text{ (Ritus, 1970 [11])}} + \underbrace{\text{[Diagram 2]}}_{\simeq \alpha^2 \chi \log \chi \text{ (Ritus, 1972 [18])}} + \underbrace{\text{[Diagram 3]}}_{\simeq \alpha^2 \chi^{2/3} \log \chi \text{ (Morozov \& Ritus, 1975 [19])}} \\
 & + \underbrace{\text{[Diagram 4]}}_{\simeq \alpha \chi^2 \log \chi \text{ (?)}} + \underbrace{\text{[Diagram 5]}}_{\simeq \alpha^3 \chi^{2/3} \log^2 \chi \text{ (Narozhny, 1979 [8])}} + \underbrace{\text{[Diagram 6]}}_{\simeq \alpha^3 \chi^{4/3} \text{ (Narozhny, 1979 [8])}} \\
 & + \underbrace{\text{[Diagram 7]}}_{\simeq \alpha^3 \chi \log^2 \chi \text{ (Narozhny, 1980 [9])}} + \underbrace{\text{[Diagram 8]}}_{\simeq \alpha^3 \chi^{5/3} \text{ (Narozhny, 1980 [9])}} + \underbrace{\text{[Diagram 9]}}_{\simeq \alpha^3 \chi^{2/3} \log^2 \chi \text{ (?)}} \\
 & + \underbrace{\text{[Diagram 10]}}_{\simeq \alpha^3 \chi^{2/3} \log^2 \chi \text{ (?)}} + \dots
 \end{aligned}$$

$\alpha^0$        $\alpha^2 \log \alpha$        $\alpha^1 \log \alpha$   
 $\alpha^1 \log \alpha$        $\alpha^2 \log^2 \alpha$        $\alpha^1$   
 $\alpha^2 \log \alpha$        $\alpha^2$        $\alpha^2 \log^2 \alpha$   
 $\alpha^2 \log^2 \alpha$        $\alpha^2 \log^2 \alpha$

Then all graphs so far suppressed by  $\alpha^{\frac{1}{2}}$  except first...(including bubble chain).

# What if $\chi = \alpha^{-2}$ ?

$$\begin{aligned}
 \frac{M}{m} = & \underbrace{\text{Diagram 1}}_{\simeq \alpha \chi^{2/3} \text{ (Ritus, 1979 [11])}} + \underbrace{\text{Diagram 2}}_{\simeq \alpha^2 \chi \log \chi \text{ (Ritus, 1972 [18])}} + \underbrace{\text{Diagram 3}}_{\simeq \alpha^2 \chi^{2/3} \log \chi \text{ (Morozov \& Ritus, 1975 [19])}} \\
 & \underbrace{\text{Diagram 4}}_{\simeq \alpha^2 \chi^{2/3} \log \chi \text{ (?)}} + \underbrace{\text{Diagram 5}}_{\simeq \alpha^3 \chi^{2/3} \log^2 \chi \text{ (Narozhny, 1979 [8])}} + \underbrace{\text{Diagram 6}}_{\simeq \alpha^3 \chi^{4/3} \text{ (Narozhny, 1979 [8])}} \\
 & \underbrace{\text{Diagram 7}}_{\simeq \alpha^3 \chi \log^2 \chi \text{ (Narozhny, 1980 [9])}} + \underbrace{\text{Diagram 8}}_{\simeq \alpha^3 \chi^{5/3} \text{ (Narozhny, 1980 [9])}} + \underbrace{\text{Diagram 9}}_{\simeq \alpha^3 \chi^{2/3} \log^2 \chi \text{ (?)}} \\
 & \underbrace{\text{Diagram 10}}_{\simeq \alpha^3 \chi^{2/3} \log^2 \chi \text{ (?)}} + \dots
 \end{aligned}$$

$\alpha^{-1/3}$  (under Diagram 2)  
 $\log \alpha$  (under Diagram 2)  
 $\alpha^{2/3} \log \alpha$  (under Diagram 3)  
 $\alpha^{2/3} \log \alpha$  (under Diagram 4)  
 $\alpha^{5/3} \log^2 \alpha$  (under Diagram 5)  
 $\alpha^{1/3}$  (under Diagram 6)  
 $\alpha \log \alpha$  (under Diagram 7)  
 $\alpha^{-1/3}$  (under Diagram 8)  
 $\alpha^{5/3} \log^2 \alpha$  (under Diagram 9)  
 $\alpha^{5/3} \log^2 \alpha$  (under Diagram 10)

Then bubble chain does grow badly, parametrically. (Consistent with behavior of resummed results shown yesterday by A. Fedotov)

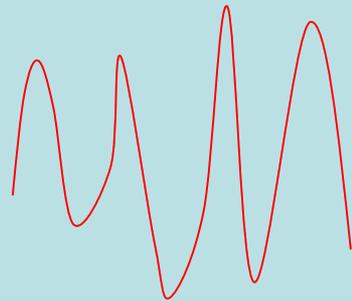
# Why might [lattice] QCD/QED help?

- Lattice gauge theory is nonperturbative.
- Successfully used to compute many quantities in hadronic physics:
  - hadron masses
  - Matrix elements between hadrons, including transition form factors
  - hadronic contributions to muon anomalous magnetic moment
- There are also at least 2 formulations of QED on the lattice (compact and noncompact).

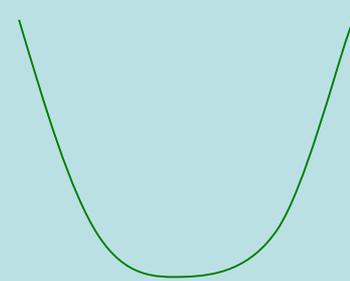
# What the lattice does very well

- Simulates gauge theories nonperturbatively in Euclidean spacetime (imaginary time):

$$Z = \int [dU(x)] e^{iS_E[U(x)]} \rightarrow \int [dU(x)] e^{-S_E[U(x)]}$$



vs.



High-dimensional integrals can be done numerically by Monte Carlo integration – if the integrands are **real** and they **don't wiggle** too much. **“Sign problem”**

# Real world is Minkowski

- Sometimes this does not matter so much, for example, computing the mass of a single, low-mass hadron:
- $$\begin{aligned}\langle O(0)O(t) \rangle &= \sum_n \langle 0|O(0)|n \rangle e^{-iM_n t} \langle n|O(t)|0 \rangle \\ &\rightarrow \sum_n \langle 0|O(0)|n \rangle e^{-M_n t} \langle n|O(t)|0 \rangle \\ &\sim e^{-M_1 t} \quad \text{as } t \rightarrow \infty\end{aligned}$$
- In Euclidean time, states decay exponentially instead of oscillating. Can determine mass of lightest hadron in a given sector reliably by waiting until heavier ones have all decayed away.

# Two hadrons

- In Minkowski spacetime, hadron-hadron scattering has complicated **imaginary phases**.
- But they affect the energy of the system in finite volume, and the **volume dependence** can be used to extract the Minkowski scattering amplitude. **Lüscher, 1986, ...**
- This is even being attempted for three hadrons, e.g. **Briceno, Hansen, Sharpe, 2016,...**
- But trying to do any more hadrons directly seems out of reach in QCD.
- Also background **electric** fields will be problematic
- Can one ask relevant “NPQED” questions for only a few quanta, and only magnetic fields?

# Thermal equilibrium

- Lattice QCD used for quark-gluon plasma studies, assuming thermal equilibrium and computing many equilibrium quantities for  $\sim 0$  baryon number

- $e^{-\beta H} \leftrightarrow$

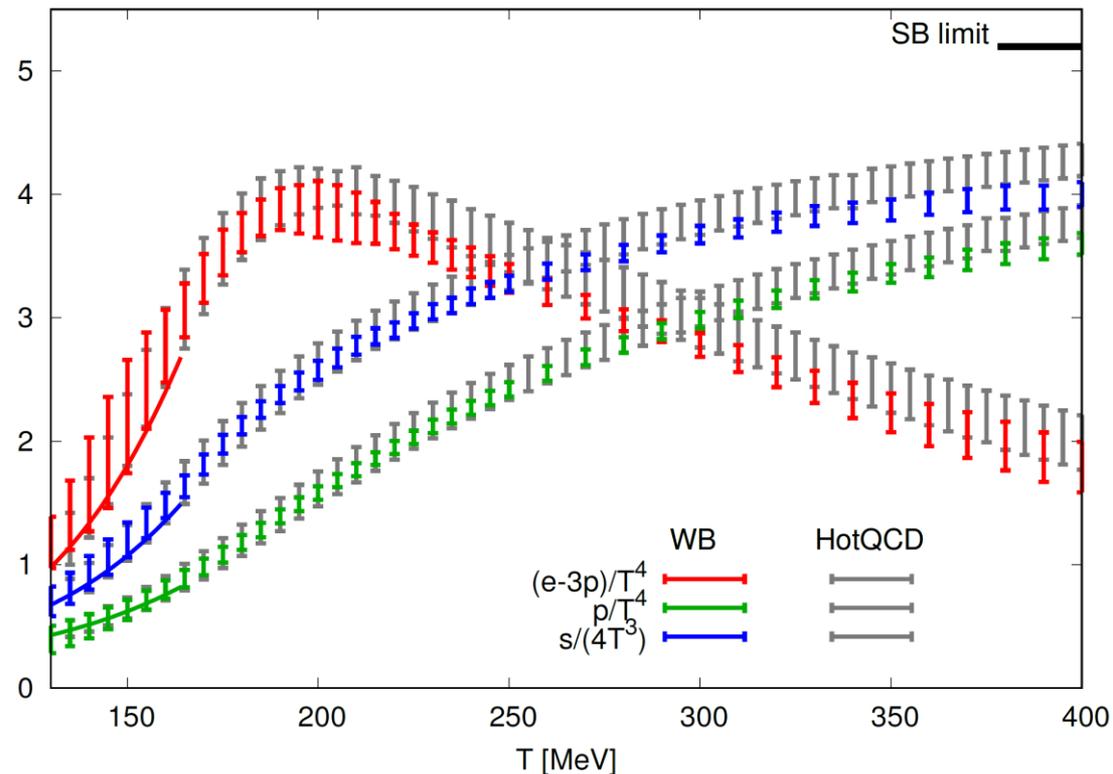
Euclidean time

periodic with period

$$\beta = 1/kT$$

Review by Ratti, 1804.07810

Lattice data for trace anomaly, entropy, pressure from HotQCD and Wuppertal Budapest (WB) collaborations



# Could similar results be useful here?

- Would have to believe some kind of equilibrium was reached, and  $\sim 0$  electron number. (Seems to be  $\sim$ OK in heavy ion collisions, despite very short time scales and many initial baryons. But equilibration times are also very short.)
- What about doing it in a large background (static?) field? Background B field probably easier (done for QCD)

# Even more exotic options for nonperturbative methods from HE theory?

- Integrability? Used to solve 1+1 dimensional Schwinger model.
- Relevant in 3+1 dim's for very high B limit [I. Shovkovy](#)
- Most applications of integrability are in 1+1, but planar N=4 super-Yang-Mills is one case in 3+1.
- Also AdS/CFT = holography has been used to gain qualitative information about other strongly coupled QFTs. How about NPQED?

Thanks for your time & patience!